

# Casimir elements for classical Lie algebras and affine Kac–Moody algebras

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# Plan of lectures

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- ▶ Affine Kac–Moody algebras: center at the critical level.
- ▶ Affine Harish-Chandra isomorphism and classical  $\mathcal{W}$ -algebras.

## Symmetric group $\mathfrak{S}_m$

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The **anti-symmetrizer** is the element

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Theorem [Jucys 1966].

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where both products are taken in the lexicographical order on the set of pairs  $(a, b)$ .

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and  $e_{ij} \in \text{End } \mathbb{C}^N$  are the matrix units.

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which we regard as elements of the algebra

$$\text{End}(\mathbb{C}^N)^{\otimes m} \cong \underbrace{\text{End} \mathbb{C}^N \otimes \dots \otimes \text{End} \mathbb{C}^N}_m.$$

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We will combine the generators into the matrix  $E = [E_{ij}]$  which will also be regarded as the element

$$E = \sum_{i,j=1}^N e_{ij} \otimes E_{ij} \in \text{End } \mathbb{C}^N \otimes U(\mathfrak{gl}_N).$$



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and for  $a = 1, \dots, m$  introduce its elements by

$$E_a = \sum_{i,j=1}^N 1^{\otimes(a-1)} \otimes e_{ij} \otimes 1^{\otimes(m-a)} \otimes E_{ij}.$$

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**Key Lemma.** The defining relations of  $U(\mathfrak{gl}_N)$  are equivalent to the single relation

$$E_1E_2 - E_2E_1 = (E_1 - E_2)P_{12}.$$

The **trace** is the linear map  $\text{End } \mathbb{C}^N \rightarrow \mathbb{C}$

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**Theorem.** For any  $s \in \mathbb{C}[\mathfrak{S}_m]$  and  $u_1, \dots, u_m \in \mathbb{C}$  the element

$$\text{tr}_{1, \dots, m} S(u_1 + E_1) \dots (u_m + E_m)$$

belongs to the center  $Z(\mathfrak{gl}_N)$  of  $U(\mathfrak{gl}_N)$ .

**Proof.** Consider the tensor product

$$\text{End } \mathbb{C}^N \otimes \text{End } (\mathbb{C}^N)^{\otimes m} \otimes U(\mathfrak{gl}_N)$$

with the copies of the algebra  $\text{End } \mathbb{C}^N$  labelled by  $0, 1, \dots, m$ .



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We will show that

$$[E_0, \text{tr}_{1, \dots, m} S(u_1 + E_1) \dots (u_m + E_m)] = 0.$$

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By the Key Lemma,

$$[E_0, u_a + E_a] = P_{0a}(u_a + E_a) - (u_a + E_a)P_{0a},$$

where we used the relations  $P_{ab}E_b = E_aP_{ab}$ .

Hence

$$\begin{aligned} & [E_0, S(u_1 + E_1) \dots (u_m + E_m)] \\ &= S \sum_{a=1}^m P_{0a} (u_1 + E_1) \dots (u_m + E_m) \\ & \quad - S(u_1 + E_1) \dots (u_m + E_m) \sum_{a=1}^m P_{0a}, \end{aligned}$$

because  $E_0 S = S E_0$  and  $P_{0a}$  commutes with  $E_b$  for  $b \neq a$ .

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The sum of the permutation operators  $P_{0a}$  commutes with  $S$  (the Schur–Weyl duality). Applying the trace  $\text{tr}_{1, \dots, m}$  and using its cyclic property we get 0. □

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Then  $C(u)$  coincides with the **column-determinant**

$$C(u) = \text{cdet} \begin{bmatrix} u + E_{11} & E_{12} & \dots & E_{1N} \\ E_{21} & u + E_{22} - 1 & \dots & E_{2N} \\ \vdots & \vdots & & \vdots \\ E_{N1} & E_{N2} & \dots & u + E_{NN} - N + 1 \end{bmatrix} .$$

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All coefficients of the polynomial  $C(u)$  are Casimir elements.



Indeed, observe that by the Key Lemma

$$\begin{aligned} \left(1 - \frac{P_{ab}}{b-a}\right)(u + E_a - a + 1)(u + E_b - b + 1) \\ = (u + E_b - b + 1)(u + E_a - a + 1)\left(1 - \frac{P_{ab}}{b-a}\right). \end{aligned}$$

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Hence, the fusion formula for  $A^{(N)}$  gives

$$A^{(N)}(u + E_1) \dots (u + E_N - N + 1) = (u + E_N - N + 1) \dots (u + E_1) A^{(N)}$$

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It remains to note that  $\text{tr}_{1, \dots, N} A^{(N)} = 1$ .

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For instance, for  $m = 2$  we get

$$\mathrm{tr}_{1,2} P_{12} E_1 E_2 = \mathrm{tr}_{1,2} E_2 P_{12} E_2 = \mathrm{tr} E^2$$

because  $\mathrm{tr}_1 P_{12} = 1$ .



# The Newton identity

Theorem [Perelomov–Popov, 1966].

We have the identity

$$1 + \sum_{m=0}^{\infty} \frac{(-1)^m \operatorname{tr} E^m}{(u - N + 1)^{m+1}} = \frac{C(u + 1)}{C(u)}.$$

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**Proof.** Verify

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Hence,

$$\begin{aligned} C(u + 1) - C(u) &= N \operatorname{tr}_{1, \dots, N} A^{(N)} (u + E_1) \dots (u + E_{N-1} - N + 2) \\ &= N \operatorname{tr}_{1, \dots, N} A^{(N)} C(u) (u + E_N - N + 1)^{-1}. \end{aligned}$$

# Harish-Chandra isomorphism

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Given an  $N$ -tuple of complex numbers  $\lambda = (\lambda_1, \dots, \lambda_N)$ , the corresponding irreducible **highest weight representation**  $L(\lambda)$  of the Lie algebra  $\mathfrak{gl}_N$  is generated by a nonzero vector  $\xi \in L(\lambda)$  (the **highest vector**) such that

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$$E_{ij} \xi = 0 \quad \text{for } 1 \leq i < j \leq N, \quad \text{and}$$

$$E_{ii} \xi = \lambda_i \xi \quad \text{for } 1 \leq i \leq N.$$

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The mapping  $z \mapsto \chi(z)$  defines an algebra isomorphism

$$\chi : Z(\mathfrak{gl}_N) \rightarrow \mathbb{C}[l_1, \dots, l_N]^{\mathfrak{S}_N}$$

known as the **Harish-Chandra isomorphism**.

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which is the projection of the  $\mathfrak{h}$ -centralizer  $U(\mathfrak{gl}_N)^{\mathfrak{h}}$  with respect to the direct sum decomposition

$$U(\mathfrak{gl}_N)^{\mathfrak{h}} = U(\mathfrak{h}) \oplus \left( U(\mathfrak{gl}_N)^{\mathfrak{h}} \cap U(\mathfrak{gl}_N)\mathfrak{n}_+ \right).$$

**Example.** Under the Harish-Chandra isomorphism we have

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This is immediate from the definition

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By the Newton formula, the Harish-Chandra images of the

Gelfand invariants are found by

$$1 + \sum_{m=0}^{\infty} \frac{(-1)^m \chi(\operatorname{tr} E^m)}{(u - N + 1)^{m+1}} = \prod_{i=1}^N \frac{u + l_i + 1}{u + l_i}.$$

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We have

$$Z(\mathfrak{g}) = \mathbb{C}[P_1, \dots, P_n],$$

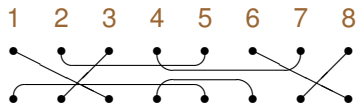
for certain algebraically independent invariants  $P_1, \dots, P_n$

whose degrees  $d_1, \dots, d_n$  are the **exponents** of  $\mathfrak{g}$  increased by 1.

Brauer algebra  $\mathcal{B}_m(\omega)$

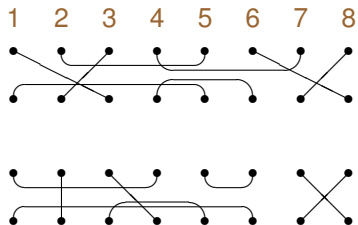
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Multiplication of  $m$ -diagrams ( $m = 8$ ):



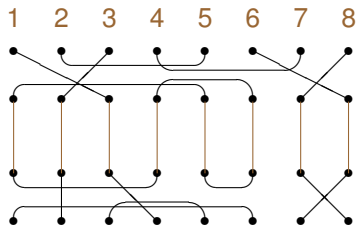
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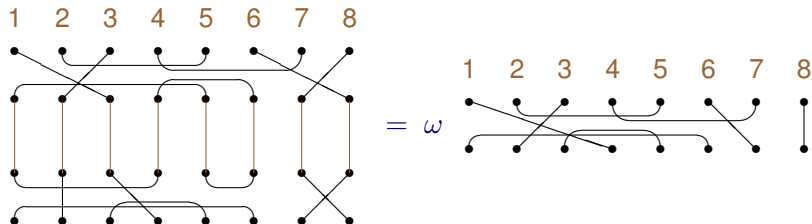
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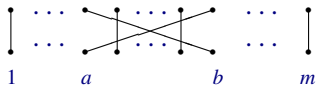




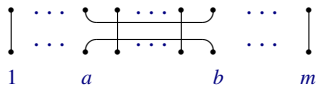
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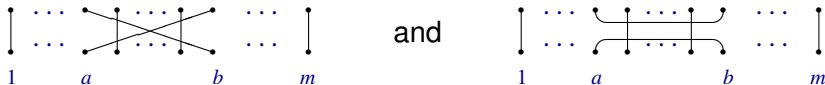


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The **symmetrizer** in  $\mathcal{B}_m(\omega)$  is the idempotent  $s^{(m)}$  such that

$$s_{ab} s^{(m)} = s^{(m)} s_{ab} = s^{(m)} \quad \text{and} \quad g_{ab} s^{(m)} = s^{(m)} g_{ab} = 0.$$

Explicitly,

$$s^{(m)} = \frac{1}{m!} \sum_{r=0}^{\lfloor m/2 \rfloor} (-1)^r \binom{\omega/2 + m - 2}{r}^{-1} \sum_{d \in \mathcal{D}^{(r)}} d,$$

where  $\mathcal{D}^{(r)} \subset \mathcal{B}_m(\omega)$  denotes the set of diagrams which have exactly  $r$  horizontal edges in the top.

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$$s^{(m)} = \prod_{1 \leq a < b \leq m} \left( 1 - \frac{g_{ab}}{\omega + a + b - 3} \right) h^{(m)},$$

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where the products are in the lexicographic order.

## Brauer–Schur–Weyl duality

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There are commuting actions of the classical groups in types  $B$ ,  $C$  or  $D$  and the Brauer algebra with a specialized parameter  $\omega$  on the tensor product space

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$$(\mathcal{B}_m(-N), Sp_N) \quad \text{with} \quad N = 2n.$$

## Action in tensors

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In the case  $\mathfrak{g} = \mathfrak{o}_N$  set  $\omega = N$ . The generators of  $\mathcal{B}_m(N)$  act in the tensor space

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where  $i' = N - i + 1$  and

$$Q_{ab} = \sum_{i,j=1}^N 1^{\otimes(a-1)} \otimes e_{ij} \otimes 1^{\otimes(b-a-1)} \otimes e_{i'j'} \otimes 1^{\otimes(m-b)}.$$

In the case  $\mathfrak{g} = \mathfrak{sp}_N$  with  $N = 2n$  set  $\omega = -N$ . The generators of  $\mathcal{B}_m(-N)$  act in the tensor space  $(\mathbb{C}^N)^{\otimes m}$  by

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In both cases denote by  $S^{(m)}$  the image of the symmetrizer  $s^{(m)}$  under the action in tensors,

$$S^{(m)} \in \underbrace{\text{End } \mathbb{C}^N \otimes \dots \otimes \text{End } \mathbb{C}^N}_m.$$



Explicitly, in the orthogonal case

$$S^{(m)} = \frac{1}{m!} \prod_{1 \leq a < b \leq m} \left( 1 + \frac{P_{ab}}{b-a} - \frac{Q_{ab}}{N/2 + b - a - 1} \right),$$

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**Remark.**  $S^{(n+1)} = 0$  for  $\mathfrak{g} = \mathfrak{sp}_{2n}$ . Consider  $\gamma_m(-2n) S^{(m)}$ ,

$$\gamma_m(\omega) = \frac{\omega + m - 2}{\omega + 2m - 2}, \quad \omega = \begin{cases} N & \text{for } \mathfrak{g} = \mathfrak{o}_N \\ -2n & \text{for } \mathfrak{g} = \mathfrak{sp}_{2n}. \end{cases}$$

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respectively.

Introduce the  $N \times N$  matrix  $F = [F_{ij}]$

$$F = \sum_{i,j=1}^N e_{ij} \otimes F_{ij} \in \text{End } \mathbb{C}^N \otimes \mathbf{U}(\mathfrak{g}).$$



**Theorem.** For any  $s \in \mathcal{B}_m(\omega)$  with  $\omega = \pm N$

and  $u_1, \dots, u_m \in \mathbb{C}$  the element

$$\mathrm{tr}_{1, \dots, m} S(u_1 + F_1) \dots (u_m + F_m)$$

belongs to the center  $Z(\mathfrak{g})$  of  $U(\mathfrak{g})$ .

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A version of the **Newton identity** also holds.

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where both sides are regarded as elements of the algebra  $\text{End } \mathbb{C}^N \otimes \text{End } \mathbb{C}^N \otimes U(\mathfrak{g})$  and

$$F_1 = \sum_{i,j=1}^N e_{ij} \otimes 1 \otimes F_{ij}, \quad F_2 = \sum_{i,j=1}^N 1 \otimes e_{ij} \otimes F_{ij}.$$

**Theorem.** For  $\mathfrak{g} = \mathfrak{o}_N$  the image of the Casimir element

$$\gamma_{2k}(N) \operatorname{tr} S^{(2k)} (F_1 - k) \dots (F_{2k} + k - 1)$$

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where  $l_i = F_{ii} + n - i + 1$  for  $i = 1, \dots, n$ .

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In particular, there is a linear basis of  $Z(\mathfrak{gl}_N)$  formed by the quantum immanants  $S_\lambda$  with  $\lambda$  running over partitions with at most  $N$  parts (Okounkov–Olshanski, 1996, 1998).

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The Harish-Chandra images  $\chi(S_\lambda)$  are the shifted Schur polynomials.

# Affine Kac–Moody algebras

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Define an invariant bilinear form on a simple Lie algebra  $\mathfrak{g}$ ,

$$\langle X, Y \rangle = \frac{1}{2h^\vee} \operatorname{tr}(\operatorname{ad} X \operatorname{ad} Y),$$

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For the classical types,  $\langle X, Y \rangle = \operatorname{const} \cdot \operatorname{tr} XY$ ,

$$h^\vee = \begin{cases} N & \text{for } \mathfrak{g} = \mathfrak{sl}_N, & \operatorname{const} = 1 \\ N - 2 & \text{for } \mathfrak{g} = \mathfrak{o}_N, & \operatorname{const} = \frac{1}{2} \\ n + 1 & \text{for } \mathfrak{g} = \mathfrak{sp}_{2n}, & \operatorname{const} = 1. \end{cases}$$

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$$[X[r], Y[s]] = [X, Y][r + s] + r \delta_{r, -s} \langle X, Y \rangle K,$$

where  $X[r] = Xt^r$  for any  $X \in \mathfrak{g}$  and  $r \in \mathbb{Z}$ .

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**Problem:** What are Casimir elements for  $\widehat{\mathfrak{g}}$ ?

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By [Kac 1974], the canonical quadratic Casimir element belongs to a **completion**  $\widetilde{U}_{-h^\vee}(\widehat{\mathfrak{g}})$  of  $U_{-h^\vee}(\widehat{\mathfrak{g}})$  with respect to the left ideals  $I_m$ ,  $m \geq 0$ , generated by  $t^m \mathfrak{g}[t]$ .

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Questions:

- ▶ Extension to Lie superalgebras.
- ▶ Extension to quantum affine algebras.

Example:  $\mathfrak{g} = \mathfrak{gl}_N$ . Defining relations for  $U(\widehat{\mathfrak{gl}}_N)$ :

$$\begin{aligned} E_{ij}[r] E_{kl}[s] - E_{kl}[s] E_{ij}[r] \\ = \delta_{kj} E_{il}[r+s] - \delta_{il} E_{kj}[r+s] + r \delta_{r,-s} \left( \delta_{kj} \delta_{il} - \frac{\delta_{ij} \delta_{kl}}{N} \right) K. \end{aligned}$$

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For all  $r \in \mathbb{Z}$  the sums

$$\sum_{i=1}^N E_{ii}[r]$$

are Casimir elements.

For  $r \in \mathbb{Z}$  set

$$C_r = \sum_{i,j=1}^N \left( \sum_{s < 0} E_{ij}[s] E_{ji}[r - s] + \sum_{s \geq 0} E_{ji}[r - s] E_{ij}[s] \right).$$

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All  $C_r$  are Casimir elements at the critical level, they belong to the **completed universal enveloping algebra**  $\tilde{U}_{-N}(\widehat{\mathfrak{gl}}_N)$ .



Introduce the (formal) Laurent series

$$E_{ij}(z) = \sum_{r \in \mathbb{Z}} E_{ij}[r] z^{-r-1}$$

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Given two Laurent series  $a(z)$  and  $b(z)$ ,

their **normally ordered product** is defined by

$$: a(z)b(z) : = a(z)_+ b(z) + b(z) a(z)_-.$$

Note

$$\sum_{r \in \mathbb{Z}} C_r z^{-r-2} = \sum_{i,j=1}^N \left( E_{ij}(z)_+ E_{ji}(z) + E_{ji}(z) E_{ij}(z)_- \right).$$

Note

$$\sum_{r \in \mathbb{Z}} C_r z^{-r-2} = \sum_{i,j=1}^N \left( E_{ij}(z) E_{ji}(z) + E_{ji}(z) E_{ij}(z) \right).$$

Hence, all coefficients of the series

$$\text{tr} : E(z)^2 : = \sum_{i,j=1}^N : E_{ij}(z) E_{ji}(z) :$$

are Casimir elements.

Similarly, all coefficients of the series

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Correction term: all coefficients of the series

$$\text{tr} : E(z)^4 : - \text{tr} : (\partial_z E(z))^2 :$$

are Casimir elements.



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Hence,  $\mathfrak{z}(\widehat{\mathfrak{g}})$  is a subalgebra of  $U(t^{-1}\mathfrak{g}[t^{-1}])$ .

## Properties:

- ▶ The subalgebra  $\mathfrak{z}(\widehat{\mathfrak{g}})$  of  $U(t^{-1}\mathfrak{g}[t^{-1}])$  is commutative.

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Any element of  $\mathfrak{z}(\widehat{\mathfrak{g}})$  is called a Segal–Sugawara vector.



Theorem (Feigin–Frenkel, 1992, Frenkel, 2007).

There exist Segal–Sugawara vectors  $S_1, \dots, S_n \in U(t^{-1}\mathfrak{g}[t^{-1}])$ ,

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Explicit constructions of such sets and a new proof of  
the theorem for the classical types  $A, B, C, D$ :

[Chervov–Talalaev, 2006, Chervov–M., 2009, M. 2013].

Example:  $\mathfrak{g} = \mathfrak{gl}_N$ .

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Set  $\tau = -d/dt$  and consider the  $N \times N$  matrix

$$\tau + E[-1] = \begin{bmatrix} \tau + E_{11}[-1] & E_{12}[-1] & \dots & E_{1N}[-1] \\ E_{21}[-1] & \tau + E_{22}[-1] & \dots & E_{2N}[-1] \\ \vdots & \vdots & \ddots & \vdots \\ E_{N1}[-1] & E_{N2}[-1] & \dots & \tau + E_{NN}[-1] \end{bmatrix}.$$

The coefficients  $\phi_1, \dots, \phi_N$  of the polynomial

$$\text{cdet}(\tau + E[-1]) = \tau^N + \phi_1 \tau^{N-1} + \dots + \phi_{N-1} \tau + \phi_N$$

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For  $N = 2$

$$\begin{aligned} \text{cdet}(\tau + E[-1]) &= (\tau + E_{11}[-1])(\tau + E_{22}[-1]) - E_{21}[-1]E_{12}[-1] \\ &= \tau^2 + \phi_1 \tau + \phi_2 \end{aligned}$$



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with

$$\phi_1 = E_{11}[-1] + E_{22}[-1],$$

$$\phi_2 = E_{11}[-1]E_{22}[-1] - E_{21}[-1]E_{12}[-1] + E_{22}[-2].$$

To get another family of Segal–Sugawara vectors, expand

$$\mathrm{tr} (\tau + E[-1])^m = \theta_{m0} \tau^m + \theta_{m1} \tau^{m-1} + \cdots + \theta_{mm}$$

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The following are Segal–Sugawara vectors for  $\mathfrak{gl}_N$ :

$$\mathrm{tr} E[-1], \quad \mathrm{tr} E[-1]^2, \quad \mathrm{tr} E[-1]^3, \quad \mathrm{tr} E[-1]^4 - \mathrm{tr} E[-2]^2.$$

The corresponding central elements in  $\tilde{U}_{-N}(\widehat{\mathfrak{gl}}_N)$  are recovered by the **state-field correspondence map**  $Y$  which takes elements of the vacuum module  $V(\mathfrak{gl}_N)$  to Laurent series in  $z$ ;

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By definition,

$$Y : E_{ij}[-1] \mapsto E_{ij}(z) = \sum_{r \in \mathbb{Z}} E_{ij}[r] z^{-r-1}.$$



Also,

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$$Y : \text{tr} E[-1]^4 - \text{tr} E[-2]^2 \mapsto \text{tr} : E(z)^4 : - \text{tr} : (\partial_z E(z))^2 :$$

Write

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**Remark.** The theorem holds in the same form for any complete set of Segal–Sugawara vectors.

Proving the Feigin–Frenkel theorem for the classical types:

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which yields a  $\mathfrak{g}[t]$ -module structure on the symmetric algebra

$$S(t^{-1}\mathfrak{g}[t^{-1}]) \cong S(\mathfrak{g}[t, t^{-1}]/\mathfrak{g}[t]).$$

Let  $X_1, \dots, X_d$  be a basis of  $\mathfrak{g}$  and let  $P = P(X_1, \dots, X_d)$  be a  $\mathfrak{g}$ -invariant in the symmetric algebra  $S(\mathfrak{g})$ .

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$$P_{(r)} = T^r P(X_1[-1], \dots, X_d[-1]), \quad r \geq 0,$$

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is a  $\mathfrak{g}[t]$ -invariant in the symmetric algebra  $S(t^{-1}\mathfrak{g}[t^{-1}])$ .

Theorem (Raïs–Tauvel 1992, Beilinson–Drinfeld 1997).

If  $P_1, \dots, P_n$  are algebraically independent generators of  $S(\mathfrak{g})^{\mathfrak{g}}$ , then the elements  $P_{1,(r)}, \dots, P_{n,(r)}$  with  $r \geq 0$  are algebraically independent generators of  $S(t^{-1}\mathfrak{g}[t^{-1}])^{\mathfrak{g}[t]}$ .

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Consider the algebra

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and recall its elements  $H^{(m)}$  and  $A^{(m)}$ .

**Theorem.** All coefficients of the polynomials in  $\tau = -d/dt$

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$$\begin{aligned}\mathrm{tr}_{1,\dots,m} A^{(m)}(\tau + E[-1]_1) \dots (\tau + E[-1]_m) \\ = \phi_{m0} \tau^m + \phi_{m1} \tau^{m-1} + \dots + \phi_{mm},\end{aligned}$$

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belong to the Feigin–Frenkel center  $\mathfrak{z}(\widehat{\mathfrak{gl}}_N)$ .

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The defining relations can be written in the form

$$\begin{aligned} E[r]_1 E[s]_2 - E[s]_2 E[r]_1 \\ = (E[r+s]_1 - E[r+s]_2) P_{12} + r \delta_{r,-s} (1 - NP_{12}). \end{aligned}$$

The required relations in the vacuum module are

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The elements  $\psi_{ma}$  and  $\theta_{ma}$  are expressed in terms of the  $\phi_{ma}$  through the **MacMahon Master Theorem** and the **Newton identities**, respectively.

The coefficients of the column-determinant are related to the  $\phi_{ma}$  through the relation

$$\text{cdet}(\tau + E[-1]) = \text{tr}_{1,\dots,N} A^{(N)} (\tau + E[-1]_1) \dots (\tau + E[-1]_N).$$

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implied by the fact that  $\tau + E[-1]$  is a **Manin matrix**.

Types *B*, *C* and *D*



## Types $B$ , $C$ and $D$

Recall the symmetrizers associated with  $\mathfrak{o}_N$  and  $\mathfrak{sp}_{2n}$ :

$$S^{(m)} = \frac{1}{m!} \prod_{1 \leq a < b \leq m} \left( 1 + \frac{P_{ab}}{b-a} - \frac{Q_{ab}}{N/2 + b - a - 1} \right),$$

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Also,

$$\gamma_m(\omega) = \frac{\omega + m - 2}{\omega + 2m - 2}, \quad \omega = \begin{cases} N & \text{for } \mathfrak{g} = \mathfrak{o}_N \\ -2n & \text{for } \mathfrak{g} = \mathfrak{sp}_{2n}. \end{cases}$$

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Combine into a matrix

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belong to the Feigin–Frenkel center  $\mathfrak{z}(\widehat{\mathfrak{g}})$ .

**Theorem.** All coefficients of the polynomial in  $\tau = -d/dt$

$$\begin{aligned} \gamma_m(\omega) \operatorname{tr}_{1, \dots, m} S^{(m)}(\tau + F[-1]_1) \dots (\tau + F[-1]_m) \\ = \phi_{m0} \tau^m + \phi_{m1} \tau^{m-1} + \dots + \phi_{mm} \end{aligned}$$

belong to the Feigin–Frenkel center  $\mathfrak{z}(\widehat{\mathfrak{g}})$ .

In addition, in the case  $\mathfrak{g} = \mathfrak{o}_{2n}$ , the **Pfaffian**

$$\operatorname{Pf} F[-1] = \frac{1}{2^n n!} \sum_{\sigma \in \mathfrak{S}_{2n}} \operatorname{sgn} \sigma \cdot F_{\sigma(1)\sigma(2)'}[-1] \dots F_{\sigma(2n-1)\sigma(2n)'}[-1]$$

belongs to  $\mathfrak{z}(\widehat{\mathfrak{o}}_{2n})$ .

Moreover,  $\phi_{22}, \phi_{44}, \dots, \phi_{2n2n}$  is a complete set of Segal–Sugawara vectors for  $\mathfrak{o}_{2n+1}$  and  $\mathfrak{sp}_{2n}$ , whereas

Moreover,  $\phi_{22}, \phi_{44}, \dots, \phi_{2n-2, 2n-2}$  is a complete set of Segal–Sugawara vectors for  $\mathfrak{o}_{2n+1}$  and  $\mathfrak{sp}_{2n}$ , whereas

$\phi_{22}, \phi_{44}, \dots, \phi_{2n-2, 2n-2}, \phi'_n$  is a complete set of Segal–Sugawara vectors for  $\mathfrak{o}_{2n}$ , where  $\phi'_n = \text{Pf } F[-1]$ .

# Affine Harish-Chandra isomorphism

For a triangular decomposition  $\mathfrak{g} = \mathfrak{n}_- \oplus \mathfrak{h} \oplus \mathfrak{n}_+$  consider the  
Harish-Chandra homomorphism

$$U(t^{-1}\mathfrak{g}[t^{-1}])^{\mathfrak{h}} \rightarrow U(t^{-1}\mathfrak{h}[t^{-1}]),$$

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the projection modulo the left ideal generated by  $t^{-1}\mathfrak{n}_-[t^{-1}]$ .

The restriction to  $\mathfrak{z}(\widehat{\mathfrak{g}})$  yields the **Harish-Chandra isomorphism**

$$f : \mathfrak{z}(\widehat{\mathfrak{g}}) \rightarrow \mathcal{W}({}^L\mathfrak{g}),$$

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where  $\mathcal{W}({}^L\mathfrak{g})$  is the **classical  $\mathcal{W}$ -algebra** associated with the Langlands dual Lie algebra  ${}^L\mathfrak{g}$  [Feigin and Frenkel, 1992].



**Example**  $\mathfrak{g} = \mathfrak{gl}_N$ . Set  $\mu_i[r] = E_{ii}[r]$ . We have

$$f : \text{cdet}(\tau + E[-1]) \mapsto (\tau + \mu_N[-1]) \dots (\tau + \mu_1[-1]).$$

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Define the elements  $\mathcal{E}_1, \dots, \mathcal{E}_N$  by the **Miura transformation**

$$(\tau + \mu_N[-1]) \dots (\tau + \mu_1[-1]) = \tau^N + \mathcal{E}_1 \tau^{N-1} + \dots + \mathcal{E}_N.$$

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Explicitly,

$$\mathcal{E}_m = e_m(T + \mu_1[-1], \dots, T + \mu_N[-1]) \mathbf{1}$$

is the **noncommutative elementary symmetric function**,

$$e_m(x_1, \dots, x_p) = \sum_{i_1 > \dots > i_m} x_{i_1} \dots x_{i_m}.$$

If  $N = 2$  then

$$\mathcal{E}_1 = \mu_1[-1] + \mu_2[-1],$$

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If  $N = 3$  then

$$\mathcal{E}_1 = \mu_1[-1] + \mu_2[-1] + \mu_3[-1],$$

$$\begin{aligned} \mathcal{E}_2 = & \mu_1[-1] \mu_2[-1] + \mu_1[-1] \mu_3[-1] + \mu_2[-1] \mu_3[-1] \\ & + 2 \mu_1[-2] + \mu_2[-2], \end{aligned}$$

$$\begin{aligned} \mathcal{E}_3 = & \mu_1[-1] \mu_2[-1] \mu_3[-1] + \mu_1[-2] \mu_2[-1] \\ & + \mu_1[-2] \mu_3[-1] + \mu_1[-1] \mu_2[-2] + 2 \mu_1[-3]. \end{aligned}$$

Then

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$$\mathcal{H}_m = h_m(T + \mu_1[-1], \dots, T + \mu_N[-1]) \mathbf{1}$$

is the **noncommutative complete symmetric function**,

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for  $N = 2n + 1$ .

For the Lie algebra  $\mathfrak{g} = \mathfrak{o}_{2n}$  the image is

$$\begin{aligned} & \frac{1}{2} h_m(\tau + \mu_1[-1], \dots, \tau + \mu_{n-1}[-1], \tau - \mu_n[-1], \dots, \tau - \mu_1[-1]) \\ & + \frac{1}{2} h_m(\tau + \mu_1[-1], \dots, \tau + \mu_n[-1], \tau - \mu_{n-1}[-1], \dots, \tau - \mu_1[-1]). \end{aligned}$$

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The Harish-Chandra image of the Pfaffian

$$\text{Pf } F[-1] = \frac{1}{2^{n_n}!} \sum_{\sigma \in \mathfrak{S}_{2n}} \text{sgn } \sigma \cdot F_{\sigma(1)\sigma(2)'}[-1] \dots F_{\sigma(2n-1)\sigma(2n)'}[-1]$$

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**Miura transformation** for  $\mathfrak{o}_{2n+1}$  [Drinfeld–Sokolov 1985]:

$$\begin{aligned} & (\tau - \mu_1[-1]) \dots (\tau - \mu_n[-1]) \tau (\tau + \mu_n[-1]) \dots (\tau + \mu_1[-1]) \\ & = \tau^{2n+1} + \mathcal{E}_2 \tau^{2n-1} + \mathcal{E}_3 \tau^{2n-2} + \dots + \mathcal{E}_{2n+1}. \end{aligned}$$

# Classical $\mathcal{W}$ -algebras

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the  $V_i$  are the **screening operators**.



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One verifies directly that

$$V_i (\tau + \mu_N[-1]) \dots (\tau + \mu_1[-1]) = 0.$$

Equivalently,

$$V_i : \mu_i(z) \mapsto \exp \int (\mu_i(z) - \mu_{i+1}(z)) dz,$$

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$$\mu_i(z) = \sum_{r=0}^{\infty} \mu_i[-r - 1] z^r, \quad i = 1, \dots, N.$$

## Affine Poisson vertex algebra $\mathcal{V}(\mathfrak{g})$

Let  $\mathfrak{g}$  be a simple Lie algebra over  $\mathbb{C}$  and

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Consider the differential algebra  $\mathcal{V} = \mathcal{V}(\mathfrak{g})$ ,

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equipped with the derivation  $\partial$ ,

$$\partial(X_i^{(r)}) = X_i^{(r+1)}$$

for all  $i = 1, \dots, d$  and  $r \geq 0$ .

Introduce the  $\lambda$ -bracket on  $\mathcal{V}$  as a linear map

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$$\{a_\lambda bc\} = \{a_\lambda b\}c + \{a_\lambda c\}b.$$

# Hamiltonian reduction



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The **classical  $\mathcal{W}$ -algebra**  $\mathcal{W}(\mathfrak{g})$  is defined by

$$\mathcal{W}(\mathfrak{g}) = \{P \in \mathcal{V}(\mathfrak{p}) \mid \rho\{X_{\lambda} P\} = 0 \text{ for all } X \in \mathfrak{n}_+\}.$$

The classical  $\mathcal{W}$ -algebra  $\mathcal{W}(\mathfrak{g})$  is a Poisson vertex algebra equipped with the  $\lambda$ -bracket

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**Motivation:** Hamiltonian equations

$$\frac{\partial u}{\partial t} = \{H \lambda u\} \Big|_{\lambda=0}$$

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De Sole, Kac and Valeri, 2013-15; Drinfeld and Sokolov, 1985.

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The  $\lambda$ -bracket (of **Virasoro–Magri**) on  $\mathcal{W}(\mathfrak{sl}_2)$  is given by

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The Hamiltonian equation with  $H = \frac{u^2}{2}$  is equivalent to the **KdV equation**

$$\frac{\partial u}{\partial t} = 3uu' - \frac{1}{2}u'''.$$

Another  $\lambda$ -bracket (of Gardner–Faddeev–Zakharov) on  $\mathcal{W}(\mathfrak{sl}_2)$

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The Hamiltonian equation with

$$K = \frac{1}{2} u^3 - \frac{1}{4} u u''$$

is also equivalent to the KdV equation.

## Generators of $\mathcal{W}(\mathfrak{gl}_N)$

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The invariant symmetric bilinear form on  $\mathfrak{gl}_N$  is defined by

$$\langle X, Y \rangle = \text{tr} XY, \quad X, Y \in \mathfrak{gl}_N.$$

Expand the determinant with entries in  $\mathcal{V}(\mathfrak{p}) \otimes \mathbb{C}[\partial]$ ,

$$\det \begin{bmatrix} \partial + E_{11} & 1 & 0 & 0 & \dots & 0 \\ E_{21} & \partial + E_{22} & 1 & 0 & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots \\ E_{N-11} & E_{N-12} & E_{N-13} & \dots & \dots & 1 \\ E_{N1} & E_{N2} & E_{N3} & \dots & \dots & \partial + E_{NN} \end{bmatrix}$$

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**Theorem.** All elements  $w_1, \dots, w_N$  belong to  $\mathcal{W}(\mathfrak{gl}_N)$ . Moreover,

$$\mathcal{W}(\mathfrak{gl}_N) = \mathbb{C}[w_1^{(r)}, \dots, w_N^{(r)} \mid r \geq 0].$$

# Chevalley-type theorem

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Let

$$\phi : \mathcal{V}(\mathfrak{p}) \rightarrow \mathcal{V}(\mathfrak{h})$$

denote the homomorphism of differential algebras defined on the generators as the projection  $\mathfrak{p} \rightarrow \mathfrak{h}$  with the kernel  $\mathfrak{n}_-$ .

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The restriction of  $\phi$  to  $\mathcal{W}(\mathfrak{g})$  is injective. The embedding

$$\phi : \mathcal{W}(\mathfrak{g}) \hookrightarrow \mathcal{V}(\mathfrak{h})$$

is often called the **Miura transformation**.

## Theorem.

The restriction of the homomorphism  $\phi$  to the classical  $\mathcal{W}$ -algebra  $\mathcal{W}(\mathfrak{g})$  yields an isomorphism

$$\phi : \mathcal{W}(\mathfrak{g}) \rightarrow \widetilde{\mathcal{W}}(\mathfrak{g}),$$

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The restriction of the homomorphism  $\phi$  to the classical  $\mathcal{W}$ -algebra  $\mathcal{W}(\mathfrak{g})$  yields an isomorphism

$$\phi : \mathcal{W}(\mathfrak{g}) \rightarrow \widetilde{\mathcal{W}}(\mathfrak{g}),$$

where  $\widetilde{\mathcal{W}}(\mathfrak{g})$  is the subalgebra of  $\mathcal{V}(\mathfrak{h})$  which consists of the elements annihilated by all screening operators  $V_i$ ,

$$\widetilde{\mathcal{W}}(\mathfrak{g}) = \bigcap_{i=1}^n \ker V_i.$$