

Symbolic dynamics of the reduced planar 3-body problem

Danya Rose, joint work with Holger Dullin



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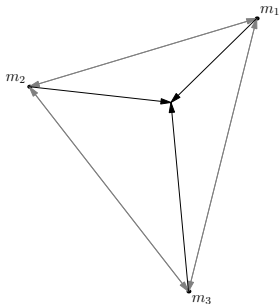


Planar 3-body problem

The problem of three massive particles moving under mutual Newtonian gravitation in 2D. Force of mass l on mass k is

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Aim: To develop a set of tools to study periodic and relatively periodic motions of the planar 3-body problem with vanishing angular momentum.

- ▶ Symmetry reduction
- ▶ Regularisation of binary collisions
- ▶ Symbolic dynamics
- ▶ Classification by discrete symmetries



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3-body Hamiltonian

Choosing units such that $G = 1$, the 3-body Hamiltonian is

$$H = T(\mathbf{P}) + U(\mathbf{X}), \text{ where}$$

$$T = \sum_{j=1}^3 \frac{P_j^2}{2m_j} \text{ and}$$

$$U = -\frac{m_2 m_3}{|X_2 - X_3|} - \frac{m_3 m_1}{|X_3 - X_1|} - \frac{m_1 m_2}{|X_1 - X_2|},$$

in which $\mathbf{X}, \mathbf{P} \in \mathbb{C}^3$ with components indexed $j = 1, 2, 3$ are, respectively, coordinates and momenta of body j and $m_j \in \mathbb{R}^+$ is its mass.

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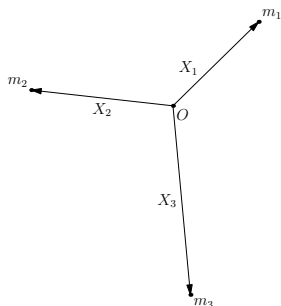
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Symmetries 1

Can remove continuous symmetries. E.g. fix centre of mass at the origin:



and extract angular momentum to separate “shape dynamics” from “rotation dynamics”.



Symmetries 2

Discrete symmetries are:

- ▶ time reversal - τ
- ▶ spatial reflection - ρ
- ▶ permutation of labels - σ_j (swap k, l , preserve j), c (cycle).

In all, forming a group of order 24 under composition.

This group is $C_2 \times C_2 \times S_3$.



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- ▶ Reduce by symmetries to remove “clutter”.
- ▶ E.g. angular momentum $L = \sum \bar{X}_j P_j$ is a constant of motion (but you wouldn't realise it from that expression right away!).
- ▶ Also reduces total dimensionality of system: 3 pairs of xy coordinates (and momenta) can be reduced to just 3 shape coordinates (and 3 momenta), plus one “rotation angle”.



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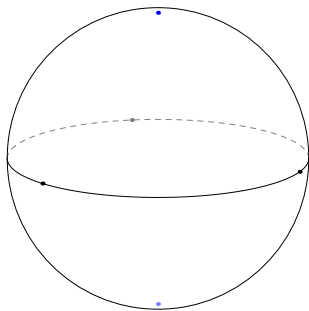


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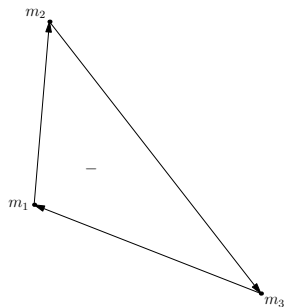
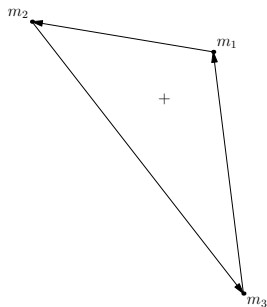
An obscure object



The shape sphere 1

Those 3 coordinates from before represent two angles and some length scale. This describes the shape space; each point is an oriented triangle.

Oriented triangles are triangles where the ordering of vertices is important:



The shape sphere 2

- ▶ Length scale is moment of inertia: $I = \sum m_j X_j$, the radial component of shape space.
- ▶ Therefore all rays from the origin represent similar oriented triangles.
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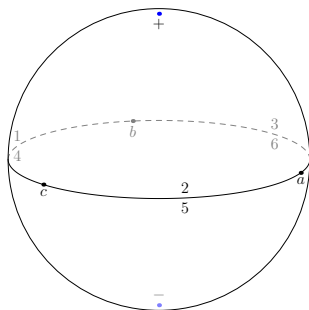


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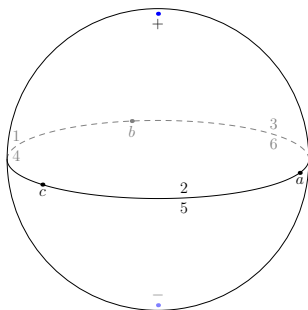
A picture of the shape sphere



- ▶ North and south poles are equilateral triangles.
- ▶ Equator is all collinear configurations.
- ▶ Three points on the equator represent binary collisions.



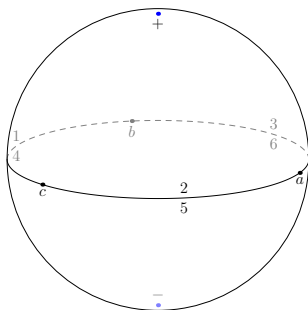
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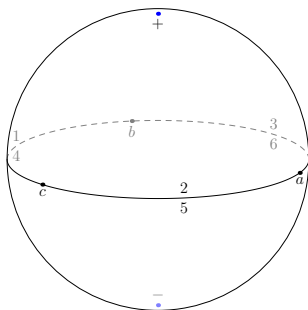
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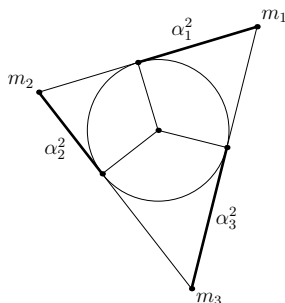
Regularising coordinates 1

Why regularise?

Even without collisions, it makes close encounters nice in numerical treatment.

Following [1], two steps to regularisation: first, make more space...

Choose new coordinates α_j s.t. $|X_l - X_k| = \alpha_k^2 + \alpha_l^2$.



- ▶ $\alpha_j = 0$ is collinearity with m_j between other two.
- ▶ $\alpha_k = \alpha_l = 0$ is a collision between m_k and m_l .



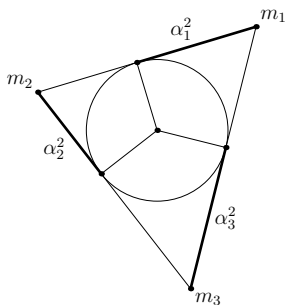
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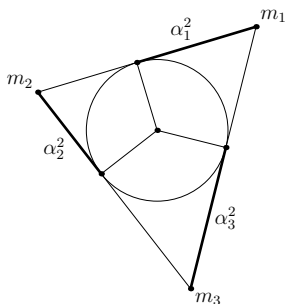
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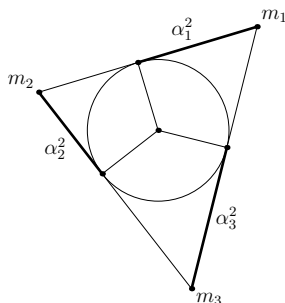
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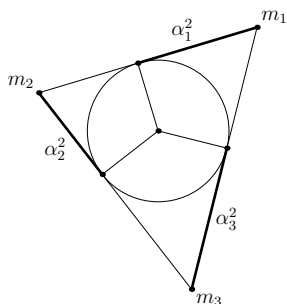
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Canonical momenta π_j are introduced, conjugate to α_j .



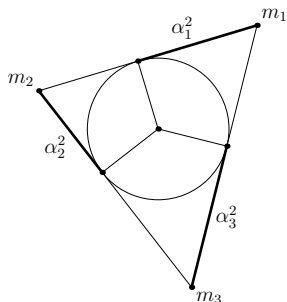
Another coordinate comes out during symmetry reduction: ϕ , a “rotation angle”.



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Scaling time

Then make more time...

Define a new time variable τ such that

$$\frac{dt}{d\tau} = g(\boldsymbol{\alpha}) = (\alpha_2^2 + \alpha_3^2)(\alpha_3^2 + \alpha_1^2)(\alpha_1^2 + \alpha_2^2).$$

Define new Hamiltonian

$$K = (H - h)g(\boldsymbol{\alpha}),$$

where H is now in new coordinates and momenta, and h is fixed to the value of H when ICs are chosen.

Physical orbits now only for $K \equiv 0$ and K is polynomial of degree 6 in $\{\alpha_j, \pi_j\}$.

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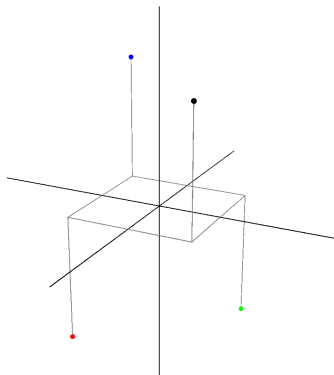
4-fold cover

New coordinates induce a 4-fold covering of shape space.

Signed area of triangle is $S = \alpha_1 \alpha_2 \alpha_3 \sqrt{\alpha_1^2 + \alpha_2^2 + \alpha_3^2}$.

Recall: $|X_l - X_k| = \alpha_k^2 + \alpha_l^2$.

So given a point $(\alpha_1, \alpha_2, \alpha_3)$, the points $(\alpha_1, -\alpha_2, -\alpha_3)$, $(-\alpha_1, \alpha_2, -\alpha_3)$ and $(-\alpha_1, -\alpha_2, \alpha_3)$ map to the same shape and orientation triangle.



New symmetry group

New Hamiltonian is

$$K = K_0 - ha_1a_2a_3, \text{ where}$$

$$K_0 = \pi^t B \pi - \sum m_k m_l a_k a_l, \text{ and}$$

$$B = \begin{pmatrix} B_1 & A_3 & A_2 \\ A_3 & B_2 & A_1 \\ A_2 & A_1 & B_3 \end{pmatrix}, \text{ with}$$

$$A_j = -\frac{a_j}{m_j} \alpha_k \alpha_l,$$

$$B_j = \frac{a_j}{m_j} \alpha^2 + \frac{a_k}{m_k} \alpha_l^2 + \frac{a_l}{m_l} \alpha_k^2 \text{ and}$$

$$\alpha^2 = \alpha_1^2 + \alpha_2^2 + \alpha_3^2,$$

which has all the original symmetries plus four new ones that act as the identity in original coordinates.



Regularised discrete symmetry group

Old symmetries act nicely on regularised coordinates, sending $(\alpha_1, \alpha_2, \alpha_3, \pi_1, \pi_2, \pi_3)$ to:

$$\tau : (\alpha_1, \alpha_2, \alpha_3, -\pi_1, -\pi_2, -\pi_3)$$

$$\rho : (-\alpha_1, -\alpha_2, -\alpha_3, -\pi_1, -\pi_2, -\pi_3)$$

$$\sigma_1 : (\alpha_1, \alpha_3, \alpha_2, \pi_1, \pi_3, \pi_2)$$

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Full group is $C_2 \times C_2 \times S_4$, order 96.



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Change of scenery: symbolic dynamics. What is it?

- ▶ Introduce an “alphabet”.
- ▶ Assign letters to certain states of the system.
- ▶ Dynamics induce a grammar.

Reduces a continuous system to a discrete one.



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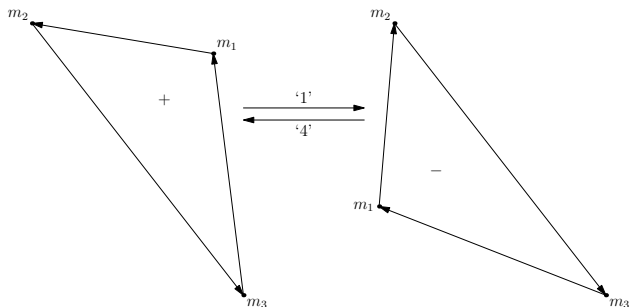
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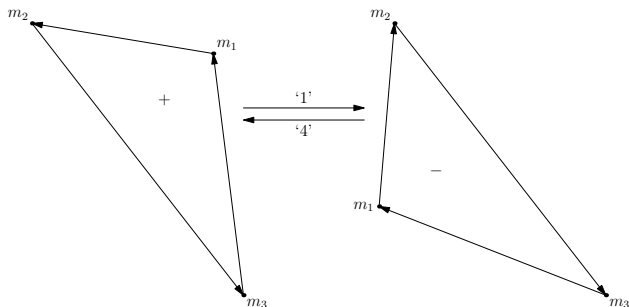
Symbolic dynamics for the 3-body problem

- ▶ Consider an orbit's projection onto the shape sphere.
- ▶ Assign symbols when passing through the equator (collinearities).
- ▶ If m_1 passes between m_2 and m_3 from upper half (positive orientation) to lower half (negative orientation) record a 1.
- ▶ Ditto from negative to positive, record a 4.
- ▶ Et cetera.



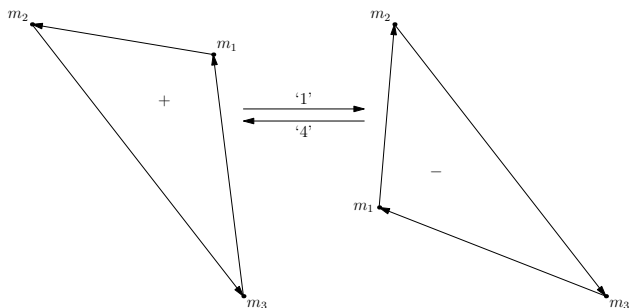
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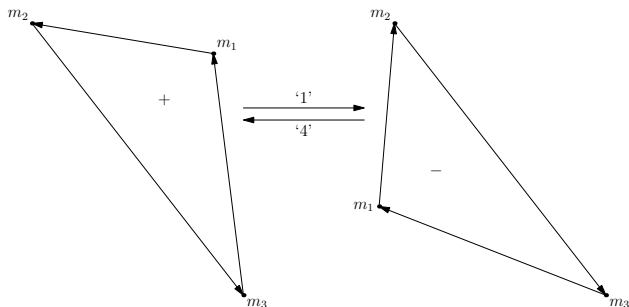
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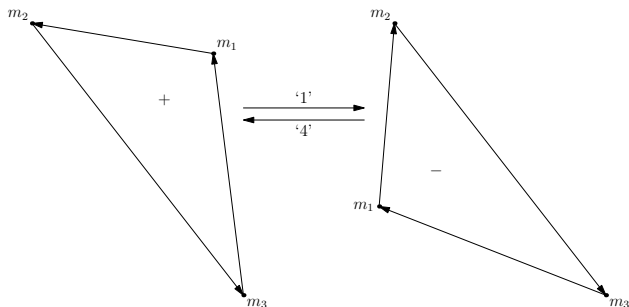
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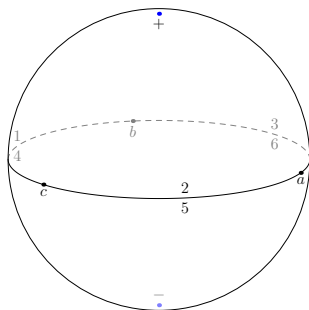
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Example: figure-8 choreography



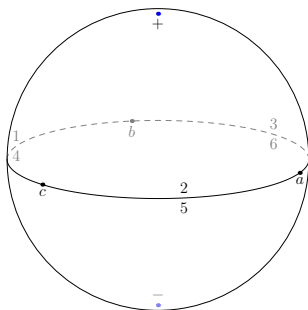
Now that we can have collisions...



- ▶ Collision points are on equator.
- ▶ Elastic collisions are orientation-preserving.
- ▶ So only touch equator, but that's still OK.



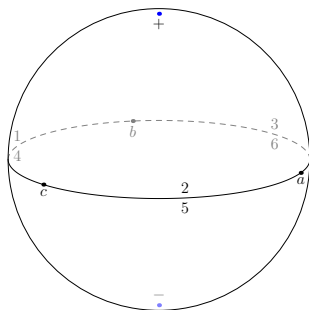
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Symmetries of the symbolic dynamics

Certain operations result in other valid symbol sequences.

Define:

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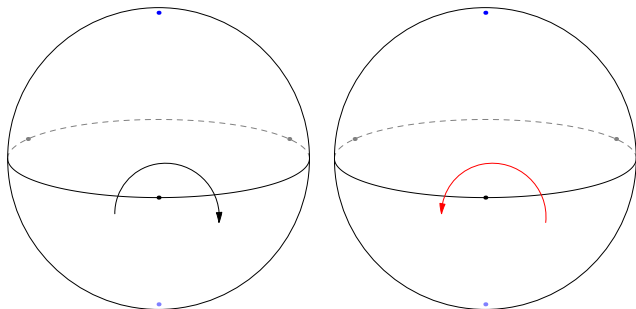
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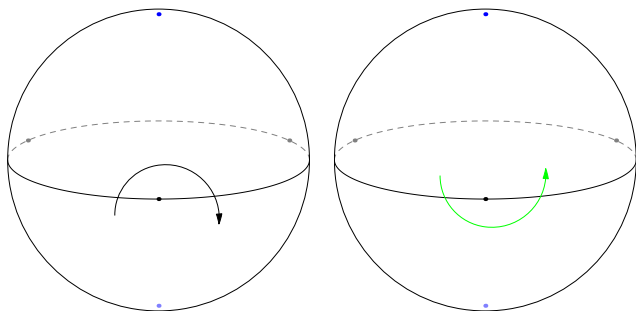


How new symmetries compare to old



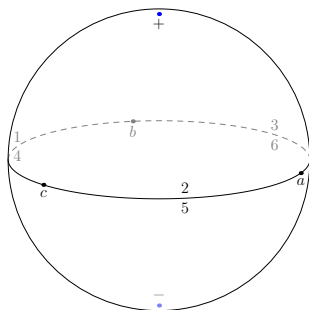
- ▶ Recall τ follows trajectory in reverse; corresponding symbol sequence is read backwards *and* orientation-reversed, corresponding to tr .
- ▶ $15 \rightarrow 24$

Symmetry comparison



- ▶ Orientation reversal ρ reflects in the equatorial plane, corresponding exactly to r .
- ▶ $15 \rightarrow 42$

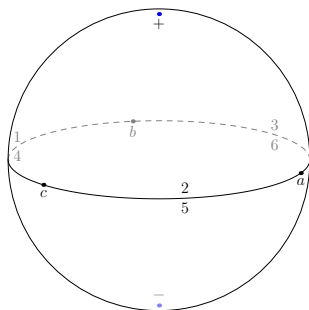
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- ▶ Index-swapping σ_j reflects in the plane through collision point of k, l and equilateral points (north/south poles). Identical between frames of reference.
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- ▶ By definition a periodic orbit repeats. If $z(\tau)$ is an orbit with period T at (scaled) time τ , then $z(\tau) = z(\tau + T)$.
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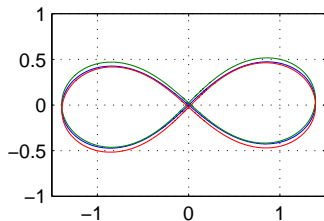
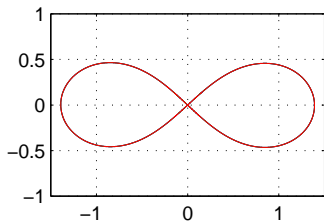
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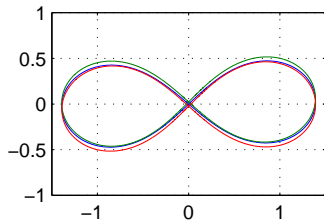
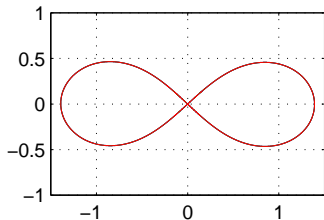
Symbol sequences are not unique!

Example: figure-8 choreography and its “shadow”.



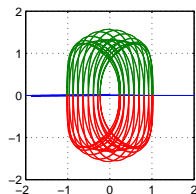
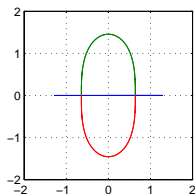
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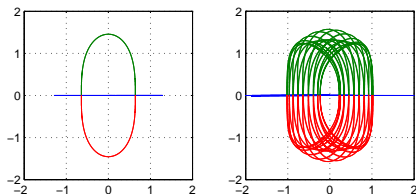
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- ▶ Orbits with same symbol sequence have same symmetries.
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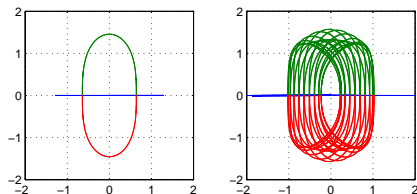
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- ▶ Orbits can be less symmetric than their symbol sequences.
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- ▶ E.g. 153426 has 12 symmetries:
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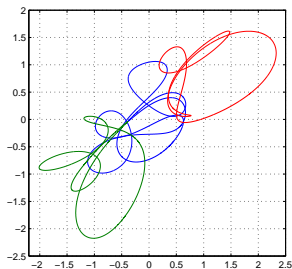
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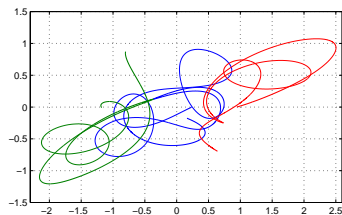
Other examples: no symmetry

Some sequences/orbits have *no* symmetries: 15151615343426



No symmetry 2

No, it actually looks like this:

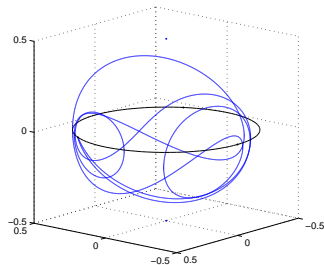


Previous picture was in a rotating frame to show that it closes.



What makes this orbit really weird

- ▶ Angle between starting and finishing configurations is -0.8262 .
- ▶ But “rotation angle” is -5.0152 .
- ▶ Reconstruction is accurate, verified by symplectic Cartesian integration with step size control.

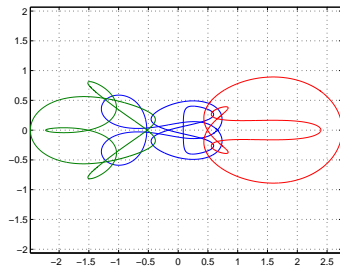


But there are well-behaved orbits

E.g. symbol sequence 15151624343426

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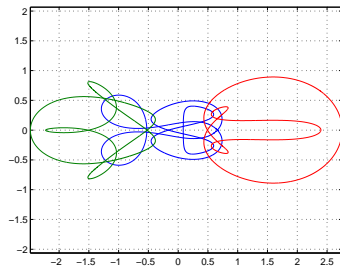


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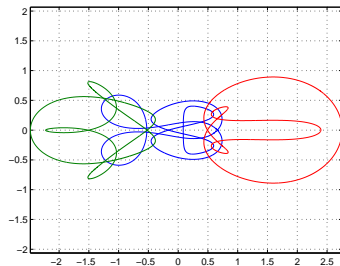


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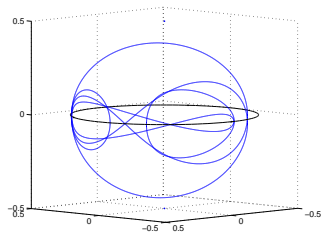
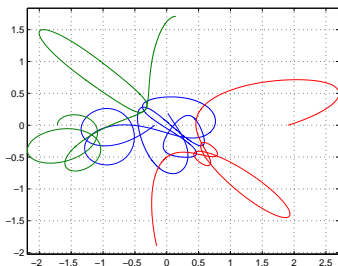
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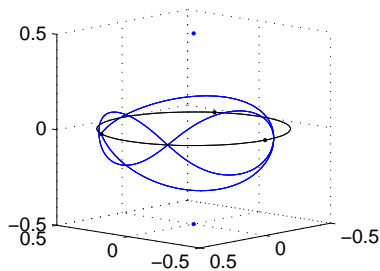
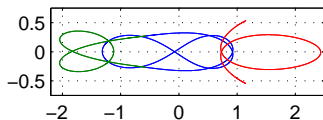
This orbit has same final Cartesian angle and rotation angle:
 -1.654 .



A collision orbit

Symbol sequence 1516c34243c6.

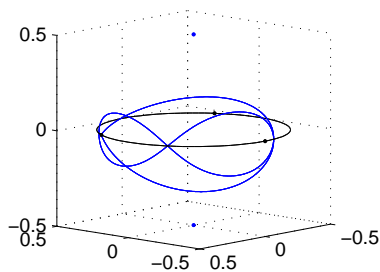
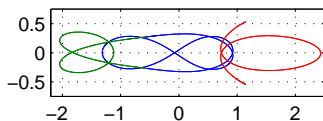
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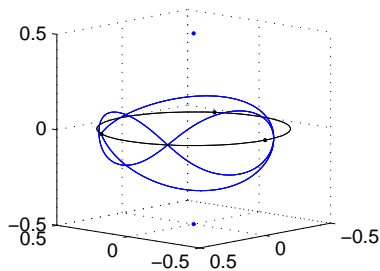
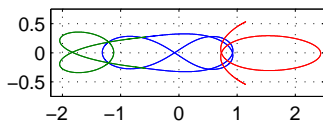
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References



Jörg Waldvogel.

Symmetric and regularized coordinates on the plane triple collision manifold.

Celestial Mechanics, 28:69–82, 1982.

